# Representation theory of a small ramified partition algebra 

P. P. MARTIN*


#### Abstract

Here $n \in \mathbb{N}, k$ is a field, $S_{n}$ is the symmetric group, $\delta \in k$, and $P_{n}(\delta)$ is the partition algebra over $k$. Our aim in this note is to study the representation theory of a subalgebra $P_{n}^{\propto}$ of $k S_{n} \otimes_{k} P_{n}(\delta)$ with certain interesting combinatorial and representation theoretic properties.

In Section 1 we discuss the motivating combinatorial background. In Section 2 we define $P_{n}^{\propto}$ (see Proposition 1). In Section 3 we determine its complex representation theory.


## 1 Introduction

The Young graph [11] has vertex set the set $\Lambda$ of all finite Young diagrams (equivalently of all integer partitions), and encodes the induction and restriction rules for ordinary irreducible modules of the sequence $\ldots \subset S_{n} \subset S_{n+1} \subset \ldots$ of symmetric groups [9]. That is, the Young graph is the Bratteli diagram for this sequence $[11, \S 1.1]$. It can be considered to lie at the heart of the analysis of these groups, and much of combinatorics [13, 11]. The multiplicity free graph (see Figure 1) and simple associated combinatorics allows a gentle build up of what, eventually, becomes a deep and powerful representation theory [8, 9, 12, 24]. In various areas of Physics [2, 17], algebra [9, 4, 25] and analysis $[1,23]$ one is led also to study the wreath products of symmetric groups:


Here however, no such multiplicity free Young graph can exist in general (at least, without further refinement), and one confronts a much more rapid onset of combinatorial complexity. By working at the level of suitable Morita equivalents,

[^0]

Figure 1: The Young graph up to rank 5
we aim to bypass this obstruction and assemble an analogue of Young's theory of comparable reach. The challenge is to find a sequence of algebras with suitable properties.

For example, as we shall see here, in the Bratteli diagram for the sequence $\ldots \subset P_{n}^{\ltimes} \subset P_{n+1}^{\ltimes} \subset \ldots$, the vertex set $\Lambda$ is replaced by $\Lambda^{\Lambda^{*}}$, the set of functions from $\Lambda^{*}$ to $\Lambda$ (where $\Lambda^{*}$ is the set of finite Young diagrams excluding the empty diagram); while the combinatorial representation theory can also be closely tied to that of wreaths. (We do not claim here that this sequence, encountered by chance while working on a different problem, is the ideal tool for this purpose, but at least that it is worth studying.) The original idea for this approach (inflating, in the sense of [14], from one combinatorial structure to another through Morita equivalences) comes serendipitously from observations on the representation theory of the ramified partition algebra [21].

Partition algebras play potentially important roles in Statistical Mechanics, in combinatorics, and in invariant theory. This is partly captured by the SchurWeyl duality diagram here:

where the groups/algebras in each layer give a dual pair of (left) actions on tensor space. In each successive layer the action shown on the left-hand side is included in the one above (and the action on the right correspondingly includes the one above). Thus $O(N)$ and $B_{n}(N)$ are the orthogonal group and the Brauer algebra respectively $[6] ; S_{N}$ acts by permuting the standard ordered basis of $k^{N}$ and $P_{n}(N)$ acts by the Potts representation [18, §8.2.1],[10]. The $S_{N-1}$ layer corresponds to breaking the global $S_{N}$ symmetry of the $N$-state Potts model by applying a magnetic field[20]. (From an invariant theory perspective this dual pair sequence has been extended below $\left(S_{N}, P_{n}(N)\right)$ in a number of ways. See for example [4, 22].)

The complex reductive representation theory (i.e. Cartan decomposition matrices and so forth, in case $k=\mathbb{C}$ ) of all the algebras appearing in this diagram is reasonably well understood. It has been noted that ramified partition algebras (RPAs) have applications in similar areas [21], but these are much less well understood. Particularly intriguing is the relationship between RPAs and wreaths (which, independently, also have roles in Physics [2, 17] and combinatorics [16]). The ramified partition algebras $P_{n}^{2}\left(\delta^{\prime}, \delta\right)$ are physically motivated subalgebras of $P_{n}\left(\delta^{\prime}\right) \otimes_{k} P_{n}(\delta)$ (see [21] for a definition; $\delta, \delta^{\prime}$ are independently chosen parameters). As we shall see in Section 2, we have algebra inclusions

$$
\begin{array}{llll}
P_{n}\left(\delta^{\prime}\right) \otimes_{k} P_{n}(\delta) & \supset k S_{n} \otimes_{k} P_{n}(\delta) & \supset & P_{n}^{\ltimes} \\
& \supset P_{n}^{2}\left(\delta^{\prime}, \delta\right) & \supset &
\end{array}
$$

and the representation theory of $P_{n}^{\ltimes}$ provides, from one perspective, a kind of approximation to that of $P_{n}^{2}$ (and hence also to that of the assembly of wreaths). Here, focussing on the representation theory of $P_{n}^{\ltimes}$, we are able to get pleasingly complete results on this representation theory (see the Theorems in the main section, §3.4). The Bratteli diagram motivating the connection to wreath combinatorics is then discussed in the final section.

## 2 Definitions

Set $\underline{n}=\{1,2, \ldots, n\}$ and $\underline{n}^{\prime}=\left\{1^{\prime}, 2^{\prime}, \ldots, n^{\prime}\right\}$ and so on. Write

$$
\operatorname{add}^{\prime}: \underline{n} \cup \underline{n}^{\prime} \rightarrow \underline{n}^{\prime} \cup \underline{n}^{\prime \prime}
$$

for the map that adds a prime; and $\operatorname{cor}^{-1}: \underline{n} \cup \underline{n}^{\prime \prime} \rightarrow \underline{n} \cup \underline{n}^{\prime}$ for the map that removes a prime when necessary (i.e. when there are two).

For $S$ a set, $\mathcal{P}_{S}$ is the set of partitions of $S$, and $\mathfrak{P}(S)$ the power set. Thus $\left|\mathcal{P}_{\underline{n}}\right|=B_{n}$, the Bell number [15]. We write ( $\mathcal{P}_{S},>$ ) for the usual refinement order on $\mathcal{P}_{S}$, that is $p>q$ if each part of $p$ is a union of parts of $q$. This order is a lattice.

Define $\mathcal{P}_{n}=\mathcal{P}_{\underline{n} \cup \underline{n}^{\prime}}$. The propagating number

$$
\#: \mathcal{P}_{n} \rightarrow \mathbb{N}
$$

is the number of parts containing both primed and unprimed elements. We write $\mathcal{P}_{n}^{\prime}$ for the subset of partitions in $\mathcal{P}_{n}$ in which every part contains both primed and unprimed elements.

The algebra $P_{n}(\delta)$ has a basis $\mathcal{P}_{n}$. We now briefly recall the algebra product. (We refer the reader to [19] or [21] for a gentler exposition.) For $a \subset \mathfrak{P}(S)$ (some $S$ ) write $\bar{a} \in \mathcal{P}_{S}$ for the most refined (lowest) partition such that each part of $\bar{a}$ is a union of elements of $a$. Thus for example $a=\{\{1,2\},\{2,3\},\{4\}\}$ gives $\bar{a}=\{\{1,2,3\},\{4\}\}$. Note that if $p, q \in \mathcal{P}_{\underline{n} \cup \underline{n}^{\prime}}$ then $p \cup \operatorname{add}^{\prime}(q) \subset \mathfrak{P}\left(\underline{n} \cup \underline{n}^{\prime} \cup \underline{n}^{\prime \prime}\right)$ and we can define

$$
p \nabla q:=\overline{p \cup \operatorname{add}^{\prime}(q)} \in \mathcal{P}_{\underline{n} \cup \underline{n}^{\prime} \cup \underline{n}^{\prime \prime}}
$$

For $r \in \mathcal{P}_{\underline{n}} \cup \underline{n}^{\prime} \cup \underline{n}^{\prime \prime}$ we write $\operatorname{res}(r)$ for the restriction of this partition to $\mathcal{P}_{\underline{n} \cup \underline{n}^{\prime \prime}}$ (so that $\left.\overline{\operatorname{cor}}^{-\prime}(\overline{\operatorname{res}}(r)) \in \mathcal{P}_{\underline{n} \cup \underline{n}^{\prime}}\right)$; and $c(r)$ for the number of parts containing only elements of $\underline{n}^{\prime}$. Then the multiplication in $P_{n}(\delta)$ is defined on pairs $p, q$ from the basis $\mathcal{P}_{n}$ by

$$
p \cdot q=\delta^{c(p \nabla q)} \operatorname{cor}^{-\prime}(\operatorname{res}(p \nabla q))
$$

Note from this construction that the set $\mathcal{P}_{n}^{\prime}$ forms a submonoid in $P_{n}(\delta)$, and that this submonoid contains an isomorphic image of $S_{n}$, defined by identifying the transposition $\sigma_{i}=(i, i+1) \in S_{n}$ with the partition

$$
\sigma_{i}=\left\{\left\{1,1^{\prime}\right\},\left\{2,2^{\prime},\right\}, \ldots,\left\{i,(i+1)^{\prime}\right\},\left\{(i+1), i^{\prime}\right\}, \ldots,\left\{n, n^{\prime}\right\}\right\}
$$

Write diag- $\mathcal{P}_{n}$ for the subset of $\mathcal{P}_{n}^{\prime}$ of partitions such that $i, i^{\prime}$ are in the same part for all $i$. Such partitions are in natural bijection with the partitions of $\underline{n}$, so $\left|\operatorname{diag}-\mathcal{P}_{n}\right|=\left|\mathcal{P}_{\underline{n}}\right|$. For example

$$
A^{i, j}:=\left\{\left\{1,1^{\prime}\right\},\left\{2,2^{\prime},\right\}, \ldots,\left\{i, i^{\prime}, j, j^{\prime}\right\}, \ldots,\left\{n, n^{\prime}\right\}\right\}
$$

is in diag- $\mathcal{P}_{n}$.
We write $M_{n}^{b}$ for the monoid generated by $\left\{A^{i j}\right\}_{i j}$ and $M_{n}^{d}$ for that generated by $S_{n} \cup\left\{A^{i j}\right\}_{i j}$. Define subalgebras of $P_{n}(\delta)$ generated by subsets: $P_{n}^{d}=$
$k\left\langle S_{n}, A^{i, j}\right\rangle_{i, j}$ and $P_{n}^{b}=k\left\langle A^{i, j}\right\rangle_{i, j}$. (Note that neither subalgebra depends on $\delta$.) These are simply the monoid algebras of the monoids above.

From the form of the partition algebra product we have
Lemma $1 P_{n}^{b}$ is a commutative algebra with basis diag- $\mathcal{P}_{n}$ of idempotents. Indeed $P_{n}^{b}$ is isomorphic (via the natural bijection) to the monoid algebra of the monoid $\left(\mathcal{P}_{\underline{n}}, \wedge\right)$, where $\wedge$ is the meet operation on $\left(\mathcal{P}_{\underline{n}},>\right)$.

Note that this remark completely determines the reductive representation theory of $P_{n}^{b}$ (as for any finite commutative monoid of idempotents).

The tensor product algebra $k S_{n} \otimes_{k} P_{n}(\delta)$ has basis $S_{n} \times \mathcal{P}_{n}$. Just as for $S_{n}[9]$ and $P_{n}^{b}$, the complex representation theory of $P_{n}(\delta)$ is well understood [20], and hence so are the tensor products [5]. We get a more challenging new algebra, however, if we proceed as follows. Define an injective map

$$
\begin{gathered}
\ltimes: S_{n} \times \operatorname{diag}-\mathcal{P}_{n} \rightarrow S_{n} \times \mathcal{P}_{n} \\
(a, b) \mapsto(a, b a)
\end{gathered}
$$

Write $P_{n}^{\ltimes}$ for the free $k$-submodule of $k S_{n} \otimes_{k} P_{n}^{d}$ with basis $\ltimes\left(S_{n} \times \operatorname{diag}-\mathcal{P}_{n}\right)$.
Proposition 1 The $k$-submodule $P_{n}^{\ltimes}$ is a subalgebra of $k S_{n} \otimes_{k} P_{n}^{d}$.
Proof: Multiplication is given by $(a, b a)(c, d c)=(a c, b a d c)$, but $b a d c=b a d a^{-1} a c$, and $b a d a^{-1} \in P_{n}^{b}$.

Proposition 2 The algebra $P_{n}^{\ltimes}$ is generated by $\left(1, A^{i j}\right)$ and $\left(\sigma_{i}, \sigma_{i}\right)($ all $i, j)$, and hence by $\left(1, A^{12}\right)$ and $\left(\sigma_{i}, \sigma_{i}\right)$.

We will write $[a, b]=\ltimes(a, b)$. Thus $[a, b][c, d]=\left[a c, b a d a^{-1}\right]$ and in particular

$$
\begin{equation*}
[a, 1][1, d]=\left[a, a d a^{-1}\right] \tag{1}
\end{equation*}
$$

Note that $A^{i j} \mapsto\left(1, A^{i j}\right)$ defines a natural injection of $P_{n}^{b}$ into $P_{n}^{\ltimes}$; and $\sigma_{i} \mapsto\left(\sigma_{i}, \sigma_{i}\right)$ a natural injection of $k S_{n}$ into $P_{n}^{\ltimes}$.

Define the set of (2-)ramified partitions

$$
\mathcal{P} \frac{2}{n}=\left\{(a, b) \mid a, b \in \mathcal{P}_{n} ; a<b\right\}
$$

From [21] this is a basis for the RPA $P_{n}^{2}\left(\delta, \delta^{\prime}\right) \subset P_{n}(\delta) \otimes_{k} P_{n}\left(\delta^{\prime}\right)$. Note also from the definition of $P_{n}^{2}\left(\delta, \delta^{\prime}\right)$ in [21] that $k S_{n} \otimes_{k} P_{n}\left(\delta^{\prime}\right)$ is not a subalgebra of $P_{n}^{2}\left(\delta, \delta^{\prime}\right)$ (for example, any non-identical pair of permutations lies outside $\mathcal{P} \frac{2}{n}$ ). However

Proposition 3 We have an algebra inclusion $P_{n}^{\ltimes} \hookrightarrow P_{n}^{2}$.
Proof: It is easy to see that elements of form $[1, b]$ and $[s, 1]$ are ramified, and these generate $P_{n}^{\ltimes}$.

Remark: We shall not make explicit use of it here, but for those comfortable with the ramified diagram calculus (see in particular [21, Fig.2]) it might well be helpful to note that the diagrams for these generators may be exemplified as follows (in case $n=5$ ):


## 3 Representation theory of $P_{n}^{\ltimes}$

### 3.1 Shapes and combinatorics

The shape of a set partition is the list of sizes of parts in non-increasing order. Thus the shape of a partition of $\underline{n}$ is an integer partition of $n$. We will write $b \Vdash \mu$ if set partition $b$ has shape $\mu$.

By convention we shall express shapes in power notation:

$$
\mu=(\underbrace{\lambda_{1}, \lambda_{1}, \ldots, \lambda_{1}}_{p_{1}}, \underbrace{\lambda_{2}, \lambda_{2}, \ldots, \lambda_{2}}_{p_{2}}, \ldots) \rightsquigarrow \quad \lambda^{p}=\left(\lambda_{1}^{p_{1}}, \lambda_{2}^{p_{2}}, \ldots\right)
$$

In particular

$$
\lambda^{p}{ }_{i}=\lambda_{i}^{p_{i}}
$$

Via this notation a shape can be considered as a pair of a strictly descending integer partition $\left(\lambda_{1}, \lambda_{2}, \ldots\right)$ and a composition $\left(p_{1}, p_{2}, \ldots\right)$ of the same length.

There is a natural action of $S_{n}$ on $\mathcal{P}_{\underline{n}}$. For each $b \in \mathcal{P}_{\underline{n}}$ define $S(b)$ as the subgroup that fixes $b$. We mention two subgroups in $S(b)$ : $S^{0}(b)$ is the group the permutes within parts: $S^{0}(b) \cong\left(S_{\lambda_{1}}\right)^{\times p_{1}} \times\left(S_{\lambda_{2}}\right)^{\times p_{2}} \times \ldots \subset S_{n}$ (in case $\left.b \Vdash \lambda^{p}\right)$; and $S^{1}(b)$ permutes parts of equal order: $S^{1}(b) \cong S_{p_{1}} \times S_{p_{2}} \times \ldots \subset S_{n}$. We have

$$
\begin{equation*}
S(b) \cong \times_{i}\left(S_{\lambda_{i}} \backslash S_{p_{i}}\right) \tag{2}
\end{equation*}
$$

Considering $S(b)$ or otherwise, the number of parts of given shape is, from [19],

$$
\begin{equation*}
\mathcal{D}_{\lambda^{p}}=\frac{n!}{\prod_{i}\left(\left(\lambda_{i}!\right)^{p_{i}} p_{i}!\right)}=\frac{n!}{\left|S\left(b \Vdash \lambda^{p}\right)\right|} \tag{3}
\end{equation*}
$$

Write $T_{b}^{L}$ (resp. $T_{b}^{R}$ ) for a traversal of the left (resp. right) cosets of $S(b)$ in $S_{n}$. I.e. $\cup_{w \in T_{b}^{L}} w S(b)$ is a partition of $S_{n}$.

### 3.2 On representations of wreaths

We shall establish later a construction of irreducible representations of our algebra $P_{n}^{\ltimes}$ directly in terms of representations of $S(b)$. Accordingly we mention these now. (However the reader may safely skip all the standard material in this section.)

Write $\Lambda$ for the set of all integer partitions including the empty partition, and $\Lambda_{n}$ for the subset of partitions of degree $n$. For $G$ a group, write $\Lambda_{\mathbb{C}}(G)$ for an index set for ordinary irreducible representations (together, in principle, with a map to explicit representations), so that $\Lambda_{\mathbb{C}}\left(S_{n}\right)=\Lambda_{n}$. (We will use the analogous notation, $\Lambda_{\mathbb{C}}(A)$, for any algebra $A$ over the complex field.) Set $s_{G}=\left|\Lambda_{\mathbb{C}}(G)\right|$ and assume there is a natural counting. For $S, T$ any sets, write $\operatorname{Hom}(S, T)$ for the set of maps $f: S \rightarrow T$. Thus an element $V$ of $\operatorname{Hom}\left(\Lambda_{\mathbb{C}}(G), \Lambda\right)$ may be expressed as an $s_{G}$-tuple ( $V_{1}, V_{2}, \ldots, V_{s_{G}}$ ) of integer partitions (a multipartition). For any $\operatorname{Hom}(S, \Lambda)$, write $\operatorname{Hom}(S, \Lambda)_{n}$ for the subset of multipartitions of total degree $n$.

The ordinary irreducible representation theory of $S(b)$ is, in effect, fairly well understood. Since $\mathbb{C}$ is a splitting field it is enough to study the wreath factors. Now see [9]. In particular we have
Theorem 1 (Cf. [9, COR.4.4.4] or [16, §1.Appendix B])

$$
\Lambda_{\mathbb{C}}\left(G \imath S_{n}\right)=\operatorname{Hom}\left(\Lambda_{\mathbb{C}}(G), \Lambda\right)_{n}
$$

The construction of irreducible $L_{V}, V \in \Lambda_{\mathbb{C}}\left(S_{l} \backslash S_{n}\right)$ is then as follows. The datum $V$ consists (see [9] or, say, [2]) of a map $V: \Lambda_{\mathbb{C}}\left(G=S_{l}\right) \rightarrow \Lambda$ such that $\sum_{i}\left|V_{i}\right|=n$. We set $v=\left(v_{1}, v_{2}, \ldots\right)=\left(\left|V_{1}\right|,\left|V_{2}\right|, \ldots\right)$ and form a traversal $T_{v}$ of the left cosets of $S_{v}$ in $S_{n}$. Let $B^{i}$ be a basis for the irreducible representation $\mathcal{S}_{i}$ in our numbering scheme for irreducible representations of $S_{l}$ (lex order of $\Lambda_{l}$, say); and $B^{V_{i}}$ a basis for the irreducible representation $\mathcal{S}\left(V_{i}\right)$ of $S_{v_{i}}$ (note that $V_{i} \vdash v_{i}$, so this is via the usual labelling scheme). Thus

$$
B^{v}=\times_{i}\left(\left(B^{i}\right)^{\times v_{i}}\right)
$$

is a basis for the irreducible representation of $S_{l}^{\times n}$ obtained from the representations $(\underbrace{\mathcal{S}_{1}, \mathcal{S}_{1}, . ., \mathcal{S}_{1}}_{v_{1} \text { copies }}, \underbrace{\mathcal{S}_{2}, \mathcal{S}_{2}, \ldots, \mathcal{S}_{2}}_{v_{2} \text { copies }}, \ldots, \mathcal{S}_{s})$ of $S_{l}$. Set

$$
B_{V}^{v}=\times_{i}\left(\left(B^{i}\right)^{\times v_{i}} \times B^{V_{i}}\right)
$$

or rather in the order

$$
B_{V}^{v}=\left(\times_{i}\left(B^{i}\right)^{\times v_{i}}\right) \times\left(\times_{i} B^{V_{i}}\right)
$$

Then $B_{V}^{v} \times T_{v}$ can be equipped with the property of basis for an (irreducible) representation of $S_{l}$ 乙 $S_{n}$.

Let $b_{1} \otimes \ldots \otimes b_{n} \otimes\left(b_{n+1} ..\right) \otimes[t]$ be an element of this basis. If $\sigma \in S_{v}, t^{\prime} \in T_{v}$, then the action of $\left(g_{1}, g_{2}, . ., g_{n} ; t^{\prime} \sigma\right)$ is given by

$$
\left(g_{1}, g_{2}, . ., g_{n} ; t^{\prime} \sigma\right) \quad b_{1} \otimes \ldots \otimes b_{v_{1}} \otimes b_{v_{1}+1} \otimes \ldots \otimes b_{n} \otimes\left(b_{n+1} . .\right) \otimes[t]
$$

$$
=g_{1} b_{\sigma^{-1}(1)} \otimes \ldots \otimes g_{n} b_{\sigma^{-1}(n)} \otimes \sigma\left(b_{n+1} . .\right) \otimes\left[t^{\prime} \sigma t\right]
$$

where $\left[t^{\prime} \sigma t\right]$ is understood as the coset representative of the coset containing this element. (See [9] for a much more detailed exposition, but) Note that the dimension of $L_{V}$ is clear:

$$
\begin{equation*}
\operatorname{dim} L_{V}=\frac{n!}{\prod_{i} v_{i}!} \prod_{i} d_{V_{i}}\left(d_{i}\right)^{v_{i}} \tag{4}
\end{equation*}
$$

where we write $d_{\lambda}$ for the dimension of the $S_{v_{i}}$ Specht module $\mathcal{S}_{\lambda}$, and $d_{i}$ for Specht dimensions for $S_{l}$ labeled using our numbering scheme.

Recall

$$
\begin{equation*}
n!=\sum_{\lambda \vdash n} d_{\lambda}^{2} \tag{5}
\end{equation*}
$$

### 3.3 Useful decompositions of $\Lambda^{\Lambda^{*}}$

The following will be useful later.
Another way to express an integer partition in an (ascending) power notation is simply as an element $\alpha$ of $\operatorname{Hom}\left(\mathbb{N}, \mathbb{N}_{0}\right)$ of finite support. The construct $\left(1^{\alpha(1)}, 2^{\alpha(2)}, \ldots\right)$ determines an integer partition in ordinary power notation on omitting all terms $i$ such that $\alpha(i)=0$ and then reversing the order of the remaining terms.
For example $\alpha:(1,2,3,4, \ldots)=(2,4,0,0, \ldots)$ becomes $\left(2^{4}, 1^{2}\right)$.
More generally, to specify a function $\mu \in \operatorname{Hom}(S, T)$, given $\underline{x}$ an ordered list of the elements of $S$, we may write $\mu: \underline{x}=\underline{y}$, meaning $\mu\left(x_{i}\right)=y_{i}$ (as in the example immediately above). But if almost all $\mu\left(x_{i}\right)=t_{0}$, with $t_{0}$ some given element of $T$, then it is convenient to write $\mu=\left(x_{i_{1}}, \mu\left(x_{i_{1}}\right)\right)\left(x_{i_{2}}, \mu\left(x_{i_{2}}\right)\right) \ldots$ where $\left\{i_{1}, i_{2}, \ldots\right\}$ is the set of $i$ such that $\mu\left(x_{i}\right) \neq t_{0}$. Depending on circumstances, the alternative layout

$$
\begin{equation*}
\mu=\frac{x_{i_{1}}}{\mu\left(x_{i_{1}}\right)} \frac{x_{i_{2}}}{\mu\left(x_{i_{2}}\right)} \ldots \tag{6}
\end{equation*}
$$

may also be useful.
In this notation our example above becomes $\alpha=\frac{2}{4} \frac{1}{2}$ (with $t_{0}=0$ ).
Let us write $\operatorname{Hom}^{f}\left(\Lambda^{*}, \Lambda\right)$ for the set of functions

$$
\mu: \Lambda^{*} \rightarrow \Lambda
$$

with only finitely many $\lambda \in \Lambda^{*}$ such that $\mu(\lambda) \neq \emptyset$. We also emphasise that $\Lambda$ is the set of integer partitions of finite integers. Thus the degree of $\mu \in$ $\operatorname{Hom}^{f}\left(\Lambda^{*}, \Lambda\right)$

$$
|\mu|=\sum_{\lambda}|\lambda \| \mu(\lambda)|
$$

is well defined. Write $\operatorname{Hom}_{N}\left(\Lambda^{*}, \Lambda\right)$ for the subset of $\operatorname{Hom}^{f}\left(\Lambda^{*}, \Lambda\right)$ of functions of degree $N$.
For example, $\operatorname{Hom}_{3}\left(\Lambda^{*}, \Lambda\right)=\left\{\frac{(3)}{(1)}, \frac{(21)}{(1)}, \frac{\left(1^{3}\right)}{(1)}, \frac{(2)}{(1)} \frac{(1)}{(1)}, \frac{\left(1^{2}\right)}{(1)} \frac{(1)}{(1)}, \frac{(1)}{(3)}, \frac{(1)}{(21)}, \frac{(1)}{\left(1^{3}\right)}\right\}$

The shape of $\mu \in \operatorname{Hom}^{f}\left(\Lambda^{*}, \Lambda\right)$ is an integer partition $\kappa(\mu)$ defined as follows. We specify via ascending power notation, in terms of which the partition is given by

$$
\alpha(i)=\sum_{\lambda \vdash i}|\mu(\lambda)|
$$

and then recast in ordinary power notation as described above.
Example: $\mu:\left(\emptyset,(1),(2),\left(1^{2}\right),(3), \ldots\right)=\left(\emptyset, \emptyset,(1),\left(1^{3}\right), \emptyset, \ldots\right)$ has $\kappa(\mu)=\left(2^{4}\right)$.
We write $\operatorname{Hom}_{\lambda^{p}}\left(\Lambda^{*}, \Lambda\right)$ for the subset of functions of shape $\lambda^{p}$. We have

$$
\operatorname{Hom}_{N}\left(\Lambda^{*}, \Lambda\right)=\bigcup_{\lambda^{p} \vdash N} \operatorname{Hom}_{\lambda^{p}}\left(\Lambda^{*}, \Lambda\right)
$$

In the simple case in which $\kappa$ has just a single 'factor' $i^{m}$ then $\operatorname{Hom}_{\left(i^{m}\right)}\left(\Lambda^{*}, \Lambda\right)$ is just the set of maps from $\Lambda_{i}$ to $\Lambda$ of total degree $m$. By Theorem 1 then,

$$
\Lambda_{\mathbb{C}}\left(S_{n} \prec S_{m}\right)=\operatorname{Hom}_{\left(n^{m}\right)}\left(\Lambda^{*}, \Lambda\right)
$$

Thus with $b \Vdash \lambda^{p}$

$$
\Lambda_{\mathbb{C}}(S(b))=\operatorname{Hom}_{\lambda^{p}}\left(\Lambda^{*}, \Lambda\right)
$$

We now have the notation to assert (as we shall show in Theorem 4)

$$
\Lambda_{\mathbb{C}}\left(P_{n}^{\ltimes}\right)=\operatorname{Hom}_{n}\left(\Lambda^{*}, \Lambda\right)
$$

### 3.4 Decomposing the regular $P_{n}^{\ltimes}$-module

We will say that the shape of $[a, b]=(a, b a)$ is the shape of $b$. It follows from (1) that the shape of $[a, b]=(a, b a)$ is unchanged by left or right multiplication by $\left[\sigma_{i}, 1\right]=\left(\sigma_{i}, \sigma_{i}\right)$.

As shapes of set partitions, integer partitions inherit a partial order from the order on set partitions themselves. E.g.

$$
\begin{align*}
\left(1^{4}\right)<\left(2,1^{2}\right) & <(3,1)<  \tag{4}\\
& <\left(2^{2}\right)<
\end{align*}
$$

Thus left or right multiplication by $\left[1, A^{i j}\right]$ either acts like 1 or takes the shape up in this order. Altogether, then, the left regular $P_{n}^{\ltimes}$-module is filtered by a poset of submodules (indeed ideals) labelled by shape. Set

$$
e_{\lambda^{p}}:=\sum_{b \Vdash \lambda^{p}}[1, b]
$$

and note that these are central elements in $P_{n}^{\ltimes}$. For example $e_{1^{n}}=[1,1]$. We have

$$
P_{n}^{\ltimes} e_{\lambda^{p}} \subset P_{n}^{\ltimes} e_{\lambda^{p^{\prime}}} \quad \Longleftrightarrow \quad \lambda^{p}>\lambda^{p^{\prime}}
$$

The sections $\mathcal{M}_{\lambda^{p}}$ of this poset each have basis the set of elements of $\ltimes\left(S_{n} \times\right.$ $\operatorname{diag}-\mathcal{P}_{n}$ ) of fixed shape. The number of basis elements of shape $\lambda^{p}$ is $n!\mathcal{D}_{\lambda^{p}}$.

We want to decompose the sections as far as possible.

As a vector space we have

$$
\begin{equation*}
\mathcal{M}_{\lambda^{p}}=\bigoplus_{b \Vdash \lambda^{p}} k\left[S_{n}, b\right]=\bigoplus_{b \Vdash \lambda^{p}} \bigoplus_{w \in T_{b}^{R}} k[S(b) w, b] \tag{7}
\end{equation*}
$$

Note that the $S(b)$-module $k\left[S_{n}, b\right]$ is isomorphic to $k S_{n}$ as an $S(b)$-module, and hence is simply $\frac{n!}{S(b) \mid}$ copies of the regular module.

Consider the quotient algebra of $P_{n}^{\ltimes}$ by all the ideals $P_{n}^{\ltimes} e_{\lambda^{p}}$ below $\lambda^{p}$. The central element $e_{\lambda^{p}}$ is idempotent in this quotient. Thus we can regard $\mathcal{M}_{\lambda^{p}}$ as an idempotent subalgebra of the quotient, with identity element $e_{\lambda^{p}}$. The category of left $\mathcal{M}_{\lambda^{p}}$-modules thus fully embeds in the category of left $P_{n}^{\ltimes}$ modules [7, §6.2], with the simple modules not hit by this embedding coming from the other $\mathcal{M}_{\lambda^{p \prime}}$.

Now consider the idempotent $\left[1, b_{0}\right], b_{0} \Vdash \lambda^{p}$, and note that in the algebra $\mathcal{M}_{\lambda^{p}}$ we have $\left[1, b_{0}\right][1, b]=\delta_{b_{0}, b}\left[1, b_{0}\right]$. We have

$$
\left[1, b_{0}\right] \mathcal{M}_{\lambda^{p}}=\left[1, b_{0}\right] \bigoplus_{w \in S_{n} ; b \Vdash \lambda^{p}}[w, b]=\bigoplus_{w \in S_{n} ; b \Vdash \lambda^{p}}\left[1, b_{0}\right][w, b]=\bigoplus_{w \in S_{n}}\left[w, b_{0}\right]
$$

Thus

$$
\left[1, b_{0}\right] \mathcal{M}_{\lambda^{p}}\left[1, b_{0}\right]=\bigoplus_{w \in S_{n}}\left[w, b_{0}\right]\left[1, b_{0}\right]=\bigoplus_{w \in S_{n}}\left[w, b_{0} w b_{0} w^{-1}\right]=\bigoplus_{w \in S\left(b_{0}\right)}\left[w, b_{0}\right] \cong k S\left(b_{0}\right)
$$

and

$$
\begin{aligned}
& \mathcal{M}_{\lambda^{p}}\left[1, b_{0}\right] \mathcal{M}_{\lambda^{p}}=\mathcal{M}_{\lambda^{p}} \bigoplus_{w \in S_{n}}\left[w, b_{0}\right]=\left(\bigoplus_{x \in S_{n} ; b \Vdash \lambda^{p}}[x, b]\right) \bigoplus_{w \in S_{n}}\left[w, b_{0}\right] \\
& =\bigoplus_{x \in S_{n} ; b \Vdash \lambda^{p}} \bigoplus_{w \in S_{n}}[x, b]\left[w, b_{0}\right]=\bigoplus_{x \in S_{n} ; b \Vdash \lambda^{p} ; w \in S_{n}}\left[x w, b x b_{0} x^{-1}\right]=\mathcal{M}_{\lambda^{p}}
\end{aligned}
$$

Thus
Theorem 2 The algebras $\mathcal{M}_{\lambda^{p}}$ and $k S\left(b_{0}\right)$ (with $b_{0} \Vdash \lambda^{p}$ ) are Morita equivalent.

Recall that $P_{n}^{\ltimes}$ has a subalgebra isomorphic to $P_{n}^{b}$. By restricting to this we see that no two sections contain any isomorphic factors. Thus each simple factor will appear in its section with multiplicity given by the dimension of its projective cover (with this dimension bounded from below, ab initio, by the dimension of the simple itself).

It also follows that

$$
\Lambda_{\mathbb{C}}\left(P_{n}^{\ltimes}\right)=\bigcup_{\lambda^{p} \vdash n} \Lambda_{\mathbb{C}}\left(\mathcal{M}_{\lambda^{p}}\right)
$$

so we have determined $\Lambda_{\mathbb{C}}\left(P_{n}^{\ltimes}\right)$ (by Theorem 2 and the results in $\S 3.2$ - equation(2) and Theorem 1). We will unpack the details shortly.

Next we compute the dimensions of these simple modules, and the overall algebra structure.

Consider the left submodule generated by an arbitrary non-zero element $\sum_{i j} c_{i j}\left[x_{i}, y_{j}\right]$ of the $\lambda^{p}$-th section, $\mathcal{M}_{\lambda^{p}}$. Choosing $l$ so that some scalar $c_{i l} \neq 0$, then in the section,

$$
\left[1, y_{l}\right] \sum_{i j} c_{i j}\left[x_{i}, y_{j}\right]=\sum_{i j} c_{i j}\left[1, y_{l}\right]\left[x_{i}, y_{j}\right]=\sum_{i j} c_{i j}\left[x_{i}, y_{l} y_{j}\right]=\sum_{i} c_{i l}\left[x_{i}, y_{l}\right]
$$

Thus this submodule itself contains a submodule generated by $\sum_{i} c_{i l}\left[x_{i}, y_{l}\right]$. Further, by (1) this submodule contains, for every partition of shape $\lambda^{p}$, an element of this form whose partition part is that partition. (These elements are of course all linearly independent.) Thus
Lemma 2 Any submodule of $\mathcal{M}_{\lambda^{p}}$ contains a non-vanishing element of form $\sum_{i} c_{i}\left[x_{i}, b\right]$, with $b \Vdash \lambda^{p}$.
How does $P_{n}^{\ltimes}=\left\langle\left[1, A^{12}\right],\left[S_{n}, 1\right]\right\rangle$ act on this element? As noted, $\left[1, A^{12}\right]$ acts as 1 or 0 . We consider the action of $\left[S_{n}, 1\right]$ in two parts: $[S(b), 1]$; and a traversal. The first part is simply a copy of $S(b) \hookrightarrow P_{n}^{\ltimes}$, so the element in Lemma 2 generates at least a simple $S(b)$-module. But since $S(b)$ fixes $b$, the $S(b)$-module generated will be spanned by elements of this form, so there will be an element of this form which generates precisely a simple $S(b)$-module. Meanwhile the action of an element $w$ of a traversal is

$$
w \sum_{i} c_{i}\left[x_{i}, b\right]=[w, 1] \sum_{i} c_{i}\left[x_{i}, b\right]=\sum_{i} c_{i}\left[w x_{i}, b^{w}\right]
$$

Note that the right hand side generates an $S\left(b^{w}\right)$-module that is isomorphic (via the natural group isomorphism) to the original $S(b)$-module. This tells us that every $P_{n}^{\ltimes}$-submodule of $\mathcal{M}_{\lambda^{p}}$ decomposes as a vector space in to summands, indexed by $b \Vdash \lambda^{p}$, the $b$-th of which is an $S(b)$-module isomorphic (via the various group isomorphisms) to all the other summands. Clearly then, in particular every simple $P_{n}^{\ltimes}$-submodule is at least a sum (as a vector space) of $\mathcal{D}_{\lambda^{p}}$ spaces each of which is an (isomorphic) simple module for $S(b)$ for the appropriate $b$.

In particular
Proposition 4 For each inequivalent simple $S(b)$-module $L_{\mu}$ (i.e. with $\mu \in$ $\operatorname{Hom}_{\lambda^{p}}\left(\Lambda^{*}, \Lambda\right)$ and $\lambda^{p}-1$ b) of dimension $m_{\mu}$ and basis $\left\{g_{i}^{\mu} x_{\mu} \mid i=1, . ., m_{\mu}\right\}$, say (see §3.2), there is a simple $P_{n}^{\ltimes}$-module $L_{\mu}^{\ltimes}$ of dimension

$$
\begin{equation*}
\operatorname{dim} L_{\mu}^{\ltimes}=m_{\mu} \mathcal{D}_{\lambda^{p}} \tag{8}
\end{equation*}
$$

and basis $\left\{\left[w g_{i}^{\mu} x_{\mu}, b^{w}\right] \mid i=1, . ., m_{\mu}, w \in T_{b}^{L}\right\}$. The modules $\left\{L_{\mu}^{\ltimes}\right\}$ are pairwise inequivalent.

Similarly,
Theorem 3 The decomposition of the b-th summand (any b) of $\mathcal{M}_{\lambda^{p}}$ itself, $S(b)[1, b]$, into a series of simple $S(b)$-modules passes to a complete decomposition of $\mathcal{M}_{\lambda^{p}}$ into a series of simple $P_{n}^{\ltimes}$-modules of this construction.

That is, every simple $P_{n}^{\ltimes}$-module arises this way (for some $\lambda^{p}$ ).
Working over $k$ such that $k S(b)$ is split semisimple for every shape (e.g. over the complex numbers), the multiplicity of $L_{\mu}$ in the $b$-th summand is $m_{\mu} \mathcal{D}_{\lambda^{p}}$, since the summand is $\mathcal{D}_{\lambda^{p}}$ copies of the regular $S(b)$-module. Thus (or by Theorem 2) each $\mathcal{M}_{\lambda^{p}}$ is semisimple (cf. [3, §1.7] for example), and hence

Theorem 4 Let $n \in \mathbb{N}$. Over $k$ as above, $P_{n}^{\ltimes}$ is split semisimple. The simple modules may be indexed by the set $\operatorname{Hom}_{n}\left(\Lambda^{*}, \Lambda\right)$. The dimensions of the simple modules are given by (8), using (2) and (4).

We have

$$
\begin{gather*}
n!\mathcal{D}_{\lambda}=\sum_{\mu}\left(m_{\mu}^{\lambda} \mathcal{D}_{\lambda}\right)^{2} \\
n!=\sum_{\mu}\left(m_{\mu}^{\lambda}\right)^{2} \mathcal{D}_{\lambda}  \tag{9}\\
\prod_{i}\left(\left(\lambda_{i}!\right)^{p_{i}} p_{i}!\right)=\sum_{\mu}\left(m_{\mu}^{\lambda^{p}}\right)^{2}
\end{gather*}
$$

This is not a trivial identity, but it is simply the $S(b)$ version of the hook dimension formula (cf. [13]). Note that the solution to this when $\mathcal{D}_{\lambda}=1$ is given $b y$ the hook dimension formula.

### 3.5 Examples and combinatorial restriction

First we unpack Theorem 4 a little. The explicit simple module index sets for the first few algebras $P_{n}^{\ltimes}(n=1,2,3,4, \ldots)$ are given in figure 3 . In the figure they appear as the vertices in the $n$-th layer of a certain directed graph. In a vertex, i.e. in an element $\mu \in \operatorname{Hom}_{n}\left(\Lambda^{*}, \Lambda\right)$, we use here the notation $\frac{\lambda^{|\mu(\lambda)|}}{\mu(\lambda)}$ for each 'factor' (i.e. each remaining $(\lambda, \mu(\lambda))$-pair after removing pairs of form $\frac{\lambda}{\emptyset}$ ). (We further omit the brackets, if there is only a single factor in $\mu$.) The slight redundancy here (the exponent on the 'numerator' - a redundant addition to our notation in (6)) facilitates some useful consistency checking in practical calculations.

Next we turn attention to restriction rules. The graph in the figure shows the Bratteli diagram of the sequence $P_{n-1}^{\ltimes} \subset P_{n}^{\ltimes}$ for $n \leq 4$. We define another directed graph $\mathcal{G}$ with vertex set $\operatorname{Hom}\left(\Lambda^{*}, \Lambda\right)$ as follows. Consider an element $\mu$. Each non-trivial factor is of the form $\frac{\lambda}{\mu(\lambda)}$ (or $\frac{\lambda^{|\mu(\lambda)|}}{\mu(\lambda)}$ in the redundant notation) as noted. We first define a linear map $M$ from $\mathbb{Z} \operatorname{Hom}\left(\Lambda^{*}, \Lambda\right)$ to itself by

$$
\begin{equation*}
M \mu=\left.\sum_{\lambda}^{\prime} \sum_{j} \frac{\lambda}{\mu(\lambda)-e_{j}} \sum_{k} \sum_{l} \frac{\lambda-e_{k}}{\mu\left(\lambda-e_{k}\right)+e_{l}} \mu\right|_{\lambda, \lambda-e_{k}} \tag{10}
\end{equation*}
$$

where the sum $\sum_{\lambda}^{\prime}$ is over partitions not mapped to $\emptyset$ by $\mu$; and all the sums involving rows of partitions are restricted to the appropriate addable or subtractable rows as usual (if $\lambda=(1)$ then the $\sum_{k}$ nominally consists in a single
summand contributing a factor with 'numerator' (1) $-e_{1}=\emptyset$ and 'denominator' undefined - this overall-undefined factor is omitted, but the term is kept); and $\left.\mu\right|_{\lambda, \lambda-e_{k}}$ means $\mu$ with the images of $\lambda, \lambda-e_{k}$ both omitted (NB, they are replaced by the explicitly given factors). We draw an edge between $\mu$ and $\mu^{\prime}$ in $\mathcal{G}$ if $\mu^{\prime}$ appears in $M \mu$ above.

The edges up to level 4 are as in figure 3. For example $M \frac{(2)}{(1)}=\frac{(2)}{\emptyset} \frac{(1)}{\emptyset+e_{1}}=\frac{(1)}{(1)}$, and $M \frac{(1)^{2}}{(2)}=\frac{(1)}{(1)} \frac{\emptyset}{-}=\frac{(1)}{(1)}$ omitting the undefined factor. A more challenging example is $\frac{(2)^{2}(1)^{2}}{(2)\left(1^{2}\right)}$. Here we have

$$
\begin{equation*}
M \mu=\frac{(2)(1)^{3}}{(1)(21)}+\frac{(2)(1)^{3}}{(1)\left(1^{3}\right)}+\frac{(2)^{2}(1)}{(2)(1)} \tag{11}
\end{equation*}
$$

One finds by direct calculation of characters that $\mathcal{G}$ coincides with the Bratteli diagram of the sequence $P_{n-1}^{\ltimes} \subset P_{n}^{\ltimes}$ for $n \leq 4$ (i.e. as shown in the figure). We conjecture that the graphs coincide.

We do not discuss a proof here, but a heuristic explanation for the form of $\mathcal{G}$ (i.e. the form of this conjecture) is as follows. A representation of the 'righthand end' of a diagram in a typical simple $P_{n}^{\ltimes}$-module is as in Figure 2. Here we shall call a collection of 'strings' with symmetrisation $\lambda$ a $\lambda$-string. Thus the last string in line in any diagram (i.e. string $n$ ) is involved in a $\lambda$-string for some $\lambda$. The action of $P_{n-1}^{\ltimes}$ on this $P_{n}^{\ltimes}$-module excludes the last string, which has the following effect. (We assume that the reader is familiar with induction and restriction rules for $S_{n}$.) Firstly it 'destroys' a $\lambda$-string into $\mu(\lambda)$, so we need to restrict $\mu(\lambda) \rightsquigarrow \sum_{j} \mu(\lambda)-e_{j}$ (cf. the first term in (10)). At the same time this creates a new $\lambda-e_{k}$-string for each suitable $k$. And for each such $k$ this extra string gives rise to an induction on $\mu\left(\lambda-e_{k}\right)$, hence $\mu\left(\lambda-e_{k}\right) \rightsquigarrow \sum_{l} \mu\left(\lambda-e_{k}\right)+e_{l}$ for each suitable $l$. (Note that the overall degree of every term produced in this way is $n-1$, as required.)

To explicitly illustrate the application of the structure Theorem we compute the dimensions of the modules in the example (11) above. For $\frac{(2)^{2}(1)^{2}}{(2)\left(1^{2}\right)}$ itself we need $\mathcal{D}_{2^{2} 1^{2}}=\frac{6!}{(2!)^{2} 2!(1!)^{2} 2!}=45, \operatorname{dim} L_{\frac{(2)}{(2)}}=\frac{2!}{2!} d_{(2)} d_{(2)}^{2}=1$, and $\operatorname{dim} L_{\frac{(1)}{\left(1^{2}\right)}}=$ $\frac{2!}{2!} d_{\left(1^{2}\right)} d_{(1)}^{2}=1$ (from the formula (4)), giving $\operatorname{dim} L_{\mu}^{\ltimes}=45$. For the summands on the right we have $\mathcal{D}_{21^{3}}=\frac{5!}{2!3!}=10$, and $\mathcal{D}_{2^{2} 1}=\frac{5!}{(2!)^{2} 2!}=15$. We have $\operatorname{dim} L_{\frac{\left(1^{3}\right)}{(21)}}=\frac{3!}{3!} d_{(21)} d_{(1)}^{3}=2$, and all the other $\operatorname{dim} L_{V}$ S are 1. Altogether then $\operatorname{dim} L_{\frac{(2)}{(1)} \frac{\left(1^{3}\right)}{(21)}}^{\ltimes}=20, \operatorname{dim} L_{\frac{(2)}{(1) \frac{\left(1^{3}\right)}{\left(1^{3}\right)}}}^{\ltimes}=10$, and $\operatorname{dim} L_{\frac{(2)^{2} \frac{(1)}{(2)}(1)}{\ltimes}}=15$. Note finally, then, that these dimensions indeed obey $45=20+10+15$.

### 3.6 Discussion

We remark that there is an established setting in which the Faa di Bruno coefficients $\mathcal{D}_{\lambda^{p}}$ appear ensemble in a way intriguingly analogous to their role in $P_{n}^{\ltimes}$. This is in the combinatorics of Bell matrices (that is, of Taylor series of composite functions) [1]. Let $g(x)=\sum_{i=1} g_{i} \frac{x^{i}}{i!}$ be a formal power series with


Figure 2: Right-hand end of a representative diagram in a simple module over $\mathbb{C}$, with the last string marked.
$g(0)=0$. The Bell matrix is the matrix whose $j^{t h}$ column contains the coefficients of the corresponding power series for $\frac{g^{j}(x)}{j!}$ (see e.g. $[1,(13.66)]$ ). This begins
$B[g]=\left(\begin{array}{ccccccc}g_{1} & 0 & & & & & \\ g_{2} & g_{1}^{2} & 0 & & & & \\ g_{3} & 3 g_{1} g_{2} & g_{1}^{3} & 0 & & & \\ g_{4} & 4 g_{1} g_{3}+3 g_{2}^{2} & 6 g_{1}^{2} g_{2} & g_{1}^{4} & 0 & & \\ g_{5} & 5 g_{1} g_{4}+10 g_{2} g_{3} & 10 g_{1}^{2} g_{3}+15 g_{1} g_{2}^{2} & 10 g_{1}^{3} g_{2} & g_{1}^{5} & 0 & \\ g_{6} & 6 g_{1} g_{5}+15 g_{2} g_{4} & 15 g_{1}^{2} g_{4}+60 g_{1} g_{2} g_{3} & 20 g_{1}^{3} g_{3} & 15 g_{1}^{4} g_{2} & g_{1}^{6} & 0 \\ \ldots & +10 g_{3}^{2} & +15 g_{2}^{3} & +45 g_{1}^{2} g_{2}^{2} & & & \\ \cdots & & & & & & \end{array}\right)$
The coefficients (within the entries) are the Faa di Bruno coefficients. The intriguing point is that these coincide with the multiplicities $\mathcal{D}_{\lambda^{p}}$ from (3) (i.e. with $\mathcal{D}_{\lambda^{p}}$ the coefficient of $\left.\prod_{i} g_{\lambda_{i}}^{p_{i}}\right)$. It would be interesting to extend this connection to the dimensions of the simple modules of $P_{n}^{\ltimes}$.

The sum of coefficients within the $n, k$-th entry in $B[g]$ is the Stirling number of the second kind [15], [18, $\S 8.3 .2]$. The sum of coefficients in the $n$-th row is $B_{n}$. Now, these last two facts give a connection with the partition algebra (using the combinatorial analysis in [19]). But a deeper connection between representation theory and this branch of analysis would be quite interesting to explore. (We include a brief report on some partial results in this direction in the Appendix.) One feature of the Bell matrix is that it gives a kind of linear realisation of the semigroup of compositions of suitable functions. Working formally over finite commutative rings one can study the subgroup of permutations of the ring realised by polynomial functions. Once again wreaths arise in this context, so it would be interesting to articulate a connection along these lines.

Refering back to our generalised Schur-Weyl duality diagram, we note that, besides $S_{N-1}$, the centralizers for various other actions of subgroups of $S_{N}$ on tensor space, including wreath subgroups, have been studied (although not, as yet, with physical motivation). Any such commutator is, by construction, an extension of the partition algebra action; and several have been explicitly cast as such, and used to study invariant theory. See for example [4, 22].


From the original physical perspective (i.e. of statistical mechanical transfer matrix algebras [18]), the fixing of a particular wreath subgroup, for example, seems somewhat arbitrary [17]. Here, as we see, we are able to treat large collections of types of wreath all together. This both removes the arbitrariness, and also provides a means to study some universal properties. In particular studying restriction rules for the algebras $P_{n-1}^{\ltimes} \subset P_{n}^{\ltimes}$ manifests a kind of universal restriction for wreaths (see Section 3.5), much closer to the beautiful restriction rules for ordinary symmetric groups $[9,24]$ than are available for specific wreaths (i.e. via $G^{\prime} \subset G$ or $S_{n-1} \subset S_{n}$ in $\left.G \imath S_{n}\right)$. However, this is pure representation theory. The question of possible physical interpretations for $P_{n}^{\ltimes}$ remains open.

## Acknowledgements.

The earlier version of this paper that has been hosted on my webpage for the last year or so (http://www.maths.leeds.ac.uk/~ppmartin/pdf/baby08.pdf) is not fun to read, so warm thanks to Robert Marsh for reading it, and for many valuable comments; and to Mark Wildon for reading the newer draft.

## Appendix: On the Bell matrix and graph walks

Again following [19], let $\left\{v_{i} \mid i \in \mathbb{N}\right\}$ be a basis for a seminfinite vector space, and define semiinfinite matrices by

$$
\begin{array}{ll}
\mathcal{N} v_{i}=i v_{i}, \quad & \mathcal{D}_{-} v_{i}=v_{i-1} \quad\left(\text { set } v_{0}=0\right) \\
& \mathcal{M}=\mathcal{N}+\mathcal{D}_{-}
\end{array}
$$

Then the first column of $\mathcal{M}^{n}$ gives the level- $n$ Stirling numbers of the second kind (whose sum is $B_{n}$ ).

Recall that $\Lambda$ is the set of all integer partitions. Define another seminfinite vector space $\mathbb{C}^{\Lambda}$ with basis $\left\{w_{\lambda} \mid \lambda \in \Lambda\right\}$. Define a matrix $\mathcal{L}_{+}$by

$$
\mathcal{L}_{+} w_{\lambda}=\sum_{i} w_{\lambda+e_{i}}
$$

where the sum is over rows of $\lambda$ such that $\lambda+e_{i} \in \Lambda$. This is the incidence matrix of the directed Young graph (see for example [11]). Thus the $\lambda$-th entry in the first column of $\mathcal{L}_{+}^{n}$ gives the number of walks of length $n$ from the empty diagram to $\lambda$, and hence the dimension of the corresponding irreducible $\mathbb{C} S_{n^{-}}$module.

Aiming to address the restriction conjecture, our task here is to come up with an analogous graph-walk construction for the Faa di Bruno coeffients.

For each integer partition define the set of associated broken partitions as follows. A broken partition $\lambda^{+}$is a partition together with a choice of a marked row from each collection of rows of equal length. (For this purpose we consider there to be precisely one row of length zero at the bottom of each partition. Thus this row is always marked, and is not combinatorially significant.) We further require that at most one marked row of $\lambda^{+}$is other than the top row among its collection of rows of equal length. For example $\left(2,1^{4}\right)$ has four associated


Figure 4: The Young matrix up to rank 5
broken partitions, while $\left(2^{2}, 1^{2}\right)$ has three. Define $\Lambda^{+}$as the set of all broken partitions.

If a broken partition has underlying partition $\lambda$ and marked rows $i, j, \ldots, k$, we may write $\lambda^{i, j, \ldots, k}$ for this broken partition. Thus $(2,2,1,1)^{1,3} \in \Lambda^{+}$.

If $\lambda^{+}$is a broken partition, and row $i$ is marked, then adding a box to this row: $\lambda \rightarrow \lambda+e_{i}$ does not give a partition unless the marked row is the highest among those of length equal to itself. Given such a composition $\lambda+e_{i}$, we define $\lambda+e_{i}^{\prime}$ as the partition that lies in the same orbit (in the usual orbit of compositions) as $\lambda+e_{i}$.

Define another semiinfinite vector space $\mathbb{C}^{\Lambda^{+}}$with basis $\left\{w_{\lambda}^{+} \mid \lambda \in \Lambda^{+}\right\}$. Define a matrix $\mathcal{L}_{+}^{+}$acting on $\mathbb{C}^{\Lambda} \otimes \mathbb{C}^{\Lambda^{+}}$by

$$
\mathcal{L}_{+}^{+} w_{\lambda}=\sum_{\lambda+} w_{\lambda+}^{+}
$$

where the sum is over broken partitions associated to $\lambda$; and

$$
\mathcal{L}_{+}^{+} w_{\lambda+}^{+}=\sum_{i} w_{\lambda+e_{i}^{\prime}}
$$

where the sum is over marked rows $i$ of $\lambda+$, but restricted to those with the property that all other marked rows of $\lambda+$ are at the top of their collection of rows of equal length. For example, if $\lambda^{+}$is $(2,2,1,1)^{1,4}=\square_{+}^{+}$then

$$
\mathcal{L}_{+}^{+} w_{\lambda+}^{+}=w_{(2,2,1,1)+e_{4}^{\prime}}=w_{(2,2,1,2)^{\prime}}=w_{(2,2,2,1)}
$$

while $\mathcal{L}_{+}^{+} w_{(2,2,1,1)^{1,3}}^{+}=w_{(3,2,1,1)}+w_{(2,2,2,1)}$. We postpone explicit analysis of these matrices to another work, but see Figure 5 and compare with $B[g]$.

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Figure 5: Illustration of the modified Young graph. The number of walks from the root to $\lambda$ are shown for each 'ordinary' layer.
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[^0]:    *Department of Pure Mathematics, University of Leeds, Leeds, LS2 9JT, UK

